

Anderson-Darling and New Weighted Cramér-von Mises Statistics

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Abstract. Anderson-Darling statistic modifies the classical empirical process in the interval $[0, 1]$ by multiplying it by a weighting function $\psi(t) = (t(1-t))^{-1/2}$. The weighting function redistributes the test sensitivity to deviations of the alternative distribution function from the hypothetical one between different subsets of $[0,1]$. In practice, the tests can be of interest with other weighting functions. There a new formulas was proposed for eigenfunctions of the Anderson-Darling statistics. Also, it was analyzed a statistic “inverse” to the Anderson-Darling statistic with the weighting function $\psi(t) = (t(1-t))^{1/2}$. The theory is based on the use of various special functions book [1]. In practice, could be useful the Cramér-von-Mises tests with the weighting functions, belonging to family $\psi(t) = t^\alpha(1-t)^\beta$, $\alpha > -1$, $\beta > -1$. The paper contains a table of distribution for statistics with different values of the degrees $\alpha > -1$ and $\beta > -1$. The table was calculated without using the statistical simulation.

Keywords: Cramér-von-Mises test, Anderson-Darling statistic, goodness-of-fit test, weighting function, eigenvalues, eigenfunctions, statistical tables.

1. Introduction: Weighted Cramér-von Mises test

One-dimensional weighted Cramér-von Mises statistic is

$$\omega_n^2 = n \int_0^1 \psi^2(t)(F_n(t) - t)^2 dt,$$

where $F_n(t)$ is the empirical distribution function based on the observations X_1, X_2, \dots, X_n of the uniformly distributed on $[0, 1]$ random variable, and $\psi(t)$ is a weighting function. The statistic (1) designed to test the hypothesis

$$H_0 : F(t) = t, \quad \text{against the alternative } H_1 : F(t) \neq t, \quad t \in [0, 1],$$

where $F(x)$ is continuous distribution function.

If the condition

$$\int_0^1 \psi^2(t) t(1-t) dt < \infty$$

is fulfilled then the statistic ω_n^2 converges in probability to

$$\omega^2 = \int_0^1 \xi^2(t) dt,$$

where $\xi(t)$, $t \in [0, 1]$, is the Gaussian process with zero mean and the covariance function

$$K_\psi(t, \tau) = \psi(t)\psi(\tau)(\min(t, \tau) - t\tau).$$

The Gaussian process $\xi(t)$ can be developed in the Karhunen-Loève series

$$\xi(t) = \sum_{i=1}^{\infty} \frac{x_k \varphi_k(t)}{\sqrt{\lambda_k}},$$

where x_k , $k = 1, 2, \dots$, are the independent standard normal random variables, and λ_k and $\varphi_k(t)$, $i = 1, 2, \dots$, are the eigenvalues and eigenfunctions of the linear operator with the kernel $K(t, \tau)$, i.e. solutions of the Fredholm integral equation

$$\varphi(t) = \lambda \int_0^1 \psi(t)\psi(\tau)(\min(t, \tau) - t\tau)\varphi(\tau)d\tau.$$

Under the contigual alternatives $H_1 : F_n(t) = t + \delta(t)/\sqrt{n}$, $n = 1, 2, \dots$, the distribution of ω^2 is the distribution of noncentral quadratic form

$$Q = \sum_{i=1}^{\infty} \frac{(x_k + \delta_k)^2}{\lambda_k}, \text{ where } \delta_k = \int_0^1 \delta(t)\varphi_k(t)dt.$$

By twice differentiation (2) respect to t , we obtain differential equation

$$h''(t) + \lambda\psi^2(t)h(t) = 0$$

with the conditions $h(0) = h(1) = 0$. Here, $h(t) = \varphi(t)/\psi(t)$.

Deheuvels and Martynov in article [3] described for $\psi(t) = t^\beta$ the follows result. Let $\{B(t) : 0 \leq t \leq 1\}$ be the Brownian bridge. Then, for each $\beta = \frac{1}{2\nu} - 1 > -1$, the Karhunen-Loeve expansions of $\{\xi(t) = t^\beta B(t) : 0 < t \leq 1\}$ is given by

$$t^\beta B(t) = \sum_{k=1}^{\infty} \frac{x_k e_{kB}(t)}{\sqrt{\lambda_{kB}}}, \quad e_{B,k}(t) = \frac{t^{\frac{1}{2\nu} - \frac{1}{2}} J_\nu \left(z_{\nu,k} t^{\frac{1}{2\nu}} \right)}{\sqrt{\nu} J_{\nu-1} (z_{\nu,k})}, \quad 0 < t \leq 1.$$

Here, $\{\omega_k : k \geq 1\}$ are i.i.d. $N(0, 1)$ random variables, and, for $k = 1, 2, \dots$, the eigenvalues are $\lambda_k = (z_{\nu,k}/2\nu)^2$, $z_{\nu,k}$, $k = 1, 2, \dots$, are zeros of the Bessel functions $J_\nu(z)$.

It is considered also the Cramér-von Mises statistic of the form

$$\omega_n^2(a, b) = n \int_0^1 t^{2a}(1-t)^{2b}(F_n(t) - t)^2 dt$$

with the weighting function $\psi(t) = t^a(1-t)^b$, $a > -1$, $b > -1$. The well known case is the Anderson-Darling statistic

$$A_n^2 = \omega_n^2(-0.5, -0.5) = n \int_0^1 \frac{(F_n(t) - t)^2}{t(1-t)} dt$$

with $a = -0.5$ and $b = -0.5$. For the statistic A_n^2 equation (3) is transformed to

$$t(1-t)h''(t) + \lambda h(t) = 0, \quad h(0) = h(1) = 0.$$

Anderson and Darling in article [2] found that their statistic has $\lambda_k = k(k+1)$ and $h_k(t) = \sqrt{t(1-t)}P'_k(2t-1)$, $k = 1, 2, \dots$, where $P_k(t)$, $k = 1, 2, \dots$, are the Legendre polynomials. The information related to the subject of this work can also be found in book [4, 5, 5, 6].

2. New formulas for eigenfunctions for the the Anderson-Darling statistic

At first, we will propose a direct method for deriving expressions for considering eigenfunctions. We will find the possible solutions of the differential equation (4) in the form

$$h_k(t) = t(1-t)(1+a_{1,k}t+a_{2,k}t^2+\dots+a_{k-1,k}t^{k-1}) \equiv t(1-t)V_k(t), \quad k = 1, 2, \dots$$

Here, $h_1(t)$ and $h_2(t)$ are understood to be $h_1(t) = t(1-t)$ and $h_2(t) = t(1-t)(1+a_{1,2}t)$. This solution satisfy to the conditions $h(0) = h(1) = 0$.

Theorem 1 *The solutions of the equation (4) can be represented for each $\lambda = \lambda_k = k(k+1)$, $k = 1, 2, \dots$, as*

$$V_k(t) = 1 + \beta_{1,k}t + \beta_{1,k}\beta_{2,k}t^2 + \beta_{1,k}\beta_{2,k}\beta_{3,k}t^3 + \dots + \beta_{1,k}\dots\beta_{k-1,k}t^{k-1}$$

or

$$V_k(t) = (((((\beta_{k-1,k}t + 1)\beta_{k-2,k}t + \dots + 1)\beta_{3,k}t + 1)\beta_{2,k}t + 1)\beta_{1,k}t + 1,$$

where $\beta_{s,k} = 1 - \lambda_k/\lambda_s = 1 - \frac{k(k+1)}{s(s+1)}$, $s = 1, 2, \dots, k-1$.

The following theorem solution can be represents another solution with using of the hypergeometric functions.

Theorem 2 *The solutions of the equation (4) can be represented for each $\lambda = \lambda_k = k(k+1)$, $k = 1, 2, \dots$, as*

$$h_k(t) = t \cdot {}_2F_1(-k, k+1; 2; t), \quad k = 1, 2, \dots$$

This result can be derived from the fact that the equation (4) is particular case of the equation for hypergeometric function ${}_2F_1$.

Theorem 3 *The following identity is valid:*

$$(1-t)V_k(t) \equiv {}_2F_1(-k, k+1; 2; t).$$

Theorem 4 *The normalized eigenfunctions of the covariance operator corresponding to the Anderson-Darling statistic can be written as:*

$$\begin{aligned} \varphi_k(t) &= 2\sqrt{k(k+1)(2k+1)}\sqrt{t(1-t)}V_k(t) \\ &= 2\sqrt{k(k+1)(2k+1)}\sqrt{\frac{t}{(1-t)}}{}_2F_1(-k, k+1; 2; t) = \\ &= 2\sqrt{\frac{2r+1}{k(k+1)}}\sqrt{t(1-t)}P'_k(2t-1), \\ & \quad k = 1, 2, \dots, \quad t \in [0, 1]. \end{aligned}$$

3. Statistic with $\psi(t) = \sqrt{t(1-t)}$

Here, we will consider the statistic Cramér-von Mises with $\psi(t) = \sqrt{t(1-t)}$. It is “inverse” for Anderson-Darling statistic. The equation (3) have the form $h''(t) + \lambda t(1-t)h(t) = 0$ with the conditions $h(0) = 0$ and $h(1) = 0$. Its solution is

$$\begin{aligned} h(t) &= C_1 \cdot {}_1F_1\left(\frac{1}{4} - \frac{\sqrt{a}}{16}, \frac{1}{2}, \frac{\sqrt{a}(-1+2x)^2}{4}\right) \exp(\sqrt{ax}(1-x)/2) \\ &+ C_2 \cdot {}_1F_1\left(\frac{3}{4} - \frac{\sqrt{a}}{16}, \frac{3}{2}, \frac{\sqrt{a}(-1+2x)^2}{4}\right) (-1+2x) \exp(\sqrt{ax}(1-x)/2). \end{aligned}$$

Here, ${}_1F_1(a; b; z)$ is the Kummer confluent hypergeometric function. With applying the conditions $h(0) = h(1) = 0$, the following equations can be derived

$$h(0) = C_1 \cdot {}_1F_1\left(\frac{1}{4} - \frac{\sqrt{a}}{16}, \frac{1}{2}, \frac{\sqrt{a}}{4}\right) - C_2 \cdot {}_1F_1\left(\frac{3}{4} - \frac{\sqrt{a}}{16}, \frac{3}{2}, \frac{\sqrt{a}}{4}\right) = 0,$$

$$h(1) = C_1 \cdot {}_1F_1\left(\frac{1}{4} - \frac{\sqrt{a}}{16}, \frac{1}{2}, \frac{\sqrt{a}}{4}\right) + C_2 \cdot {}_1F_1\left(\frac{3}{4} - \frac{\sqrt{a}}{16}, \frac{3}{2}, \frac{\sqrt{a}}{4}\right) = 0.$$

Hence, the equation for eigenvalues is the determinant of the previous equation

$${}_1F_1\left(\frac{1}{4} - \frac{\sqrt{a}}{16}, \frac{1}{2}, \frac{\sqrt{a}}{4}\right) \cdot {}_1F_1\left(\frac{3}{4} - \frac{\sqrt{a}}{16}, \frac{3}{2}, \frac{\sqrt{a}}{4}\right) = 0.$$

This equation can be resolved by numerical methods. First zeros of the left multiplier are:

45.24420524 514.3565172 1495.598597 2988.879509 4994.177419
7511.484585 10540.79747 14082.11420 18135.43367 22700.75517

First zeros of the right multiplier are:

215.7714658 940.9700448 2178.236123 3927.526974 6188.830121
8962.140456 12247.45544 16044.77364 20354.09420 25175.41653

4. Another weighted statistics

The equation (3) with $\psi(t) = (t(1-t))^{3/2}$ has the solution with Heun triconfluent function. The same equation with $\psi^2(t) = 1 - \cos(\pi t)$ has the solution with Mathieu function.

5. Table of the distribution for statistics with $\psi(t) = (t(1-t))^\alpha$

In the Table 1 we present the quantiles of the $\omega^2(\alpha, \beta)$ distribution with $\alpha = \beta$. More detailed tables can be found in article [7].

Table 1

Quantiles of $\omega^2(\alpha, \alpha)/S$

p/α	- 0.70	- 0.50	- 0.25
0.90	4.3776	1.9330	0.7945
0.95	5.4914	2.4924	1.0437
0.99	8.2378	3.8781	1.6606
0.995	9.4680	4.4982	1.9362
S	1	1	1

0	0.5	1	2
3.4731	0.7213	1.5834	0.8245
4.6136	0.9709	2.1479	1.1284
7.4346	1.5872	3.5397	1.8769
8.6939	1.8620	4.1603	2.2105
10	10	100	1000

Acknowledgments

The work was carried out at IITP RAS and supported by Russian Science Foundation (grant RSF No. 14-50-00150).

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